

Qualitative properties of magnetic fields in scalar field cosmology

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We study the qualitative properties of the class of spatially homogeneous Bianchi type- VI_0 cosmological models containing a perfect fluid with a linear equation of state, a scalar field with an exponential potential and a uniform cosmic magnetic field, using dynamical systems techniques. We find that all models evolve away from an expanding massless scalar field model in which the matter and the magnetic field are negligible dynamically. We also find that for a particular range of parameter values the models evolve towards the usual power-law inflationary model (with no magnetic field) and, furthermore, we conclude that inflation is not fundamentally affected by the presence of a uniform primordial magnetic field. We investigate the physical properties of the Bianchi type-I magnetic field models in some detail.

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I. INTRODUCTION

There are many observations that imply the existence of magnetic fields in the Universe, and recent observations have produced firmer estimates of the strength of magnetic fields in interstellar and intergalactic space [1–3]. Current observations place an upper bound on the strength of a cosmic magnetic field, but do not conclusively determine whether such a field exists [4,5]. Although cosmologists have investigated the possible existence of a homogeneous intergalactic magnetic field of primordial origin both from a theoretical and an observational point of view for many years, research on cosmological magnetic fields has been rather marginal despite their potential importance. The reasons could be the perceived weakness of the field effects or the lack, as yet, of a consistent theory explaining the origin of cosmic magnetism. However, this situation has changed considerably recently (cf. Ref. [6]).

Primordial magnetic fields introduce new ingredients into the standard picture of the early Universe. Such a field would affect the temperature distribution of the microwave background radiation, primeval nucleosynthesis and galaxy formation. But its most direct observational effect is the Faraday rotation it would cause in linearly polarized radiation from observed extragalactic radio sources. Fundamental properties of magnetic fields include their vectorial nature, which inevitably couples the field to the spacetime geometry, and the resulting tension (i.e., the negative pressure) exerted along the field's lines of force [7,8]. The implications of such an interaction are both kinematical and dynamical; kinematically, the magnetocurvature effect tends to accelerate positively curved perturbed regions, while it decelerates regions with negative local curvature and, dynamically, the most important magnetocurvature effect is that it can reverse the pure magnetic effect on density perturbations [9–12].

In a recent analysis [8] the kinematics were considered and, assuming a spacetime filled with a perfectly conducting

barotropic fluid permeated by a weak primordial magnetic field while treating the energy density and the anisotropic pressure of the field as first-order perturbations upon the Friedmann-Lemaître-Robertson-Walker (FLRW) background, it was shown how a cosmological magnetic field can modify the expansion rate of an almost-FLRW universe. The effects of the interaction between magnetism and geometry in cosmology are subtle, and it was argued that the magnetocurvature coupling could make the field into a key player irrespective of the magnetic strength. Indeed, even weak magnetic fields can lead to appreciable effects, provided that there is a strong curvature contribution. This was illustrated by studying spatially open cosmological models containing “matter” with negative pressure, and it was found that the phase of accelerated expansion, which otherwise would have been inevitable, may not even happen. This leads to the question of the *efficiency of inflationary models in the presence of primordial magnetism*. Recall that an initial curvature era was never considered as a problem for inflation, given the smoothing power of the accelerated expansion. However, this may not be the case when a magnetic field is present, no matter how weak the latter is. A more recent work [13] has further studied the vector nature of magnetic fields and the unique coupling between magnetism and spacetime curvature in general relativity which gives rise to a variety of effects with important implications was discussed.

Observations of the high degree of isotropy of the cosmic microwave background indicate that the Universe is “almost” isotropic and spatially homogeneous (at least since the time of the last scattering). Theoretical support for this belief comes from the so-called Ehlers-Geren-Sachs [14] theorem. Reference [15] has proven a generalization of this theorem and has consequently shown that any strong magnetic fields in the Universe are ruled out. This theoretical result is model independent and includes the case of inhomogeneous magnetic fields. In further work [16] numerical constraints are placed on all types of primordial and protogalactic magnetic fields in the Universe from cosmic microwave background data.

We shall study open spatially homogeneous cosmological models containing both a uniform magnetic field and a scalar field here, partially in an attempt to address some of the

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questions raised above. Scalar field cosmological models are of great importance in the study of the early Universe. Models with a variety of self-interaction potentials have been studied, and one potential that is commonly investigated and which arises in a number of physical situations has an exponential dependence on the scalar field [17,18]. There have been a number of studies of spatially homogeneous scalar field cosmological models with an exponential potential, with particular emphasis on the possible existence of inflation in such models (Refs. [19,32,20], and references within).

Dynamical systems methods for analyzing the qualitative properties of cosmological models have proven very useful [21,22]. A qualitative analysis of cosmological models with matter and a uniform magnetic field has been presented previously [23,24]. A universe with a primordial magnetic field is necessarily anisotropic. Thus, in order to investigate the influence of the magnetic field on the dynamics of the universe the Einstein field equations must be analyzed in models more general than the FLRW models. The simple classes of the Bianchi I and Kantowski-Sachs models were discussed (cf. Ref. [25]), and Ref. [23] applied techniques from the theory of planar dynamical systems to prove qualitative results about the evolution of the class of axisymmetric Bianchi I cosmologies with matter and a primordial magnetic field.

In Ref. [24] the Einstein-Maxwell field equations for orthogonal Bianchi type- VI_0 cosmologies with a γ -law perfect fluid and a pure, homogeneous sourcefree magnetic field were written as an autonomous differential equation in terms of expansion-normalized variables. It was shown that the physical region of state space is compact, and that the differential equation admits certain invariant sets and monotone functions which play an important role in the analysis. A complete analysis of the stability properties of the equilibrium points of the differential equation and a description of the bifurcations that occur as the equation of state parameter γ varies was given. The associated dynamical system was studied and the past, intermediate and future evolution of these models was determined. All asymptotic states of the models, and the likelihood that they will occur, were described. In particular, oscillatory behavior also occurs in cosmological models with a magnetic field (and in Einstein-Yang-Mills theory in general). Further work on spatially homogeneous models with a magnetic field and a nontilted perfect fluid has been carried out recently [26,27]. In particular, Ref. [28] has generalized (to the nonpolarized solutions) the work of Ref. [24], and rigorously shown that the evolution toward the singularity is oscillatory in Bianchi type- VI_0 vacuum models.

We shall concentrate on the Bianchi type- VI_0 models primarily for mathematical simplicity. It is well known [29] that a pure magnetic field is only possible in Bianchi cosmologies of types I, II, VI_0 , VII_0 (in class A) and type III (in class B). For types VI_0 and VII_0 , the algebraic constraints that arise from the Einstein field equations imply that the shear eigenframe is Fermi propagated, which in turn implies that the remaining field equations reduce to an autonomous differential equation with a polynomial vector field. Finally, for type VI_0 but not for VII_0 , the physical region of state space is

compact. We note that since a magnetic field is not compatible with Bianchi types VIII and IX (for example), the class of Bianchi type- VI_0 magnetic cosmologies is of the same generality as the Bianchi type-VIII/IX models (without magnetic field) [30].

In this paper we shall investigate the class of Bianchi VI_0 models with barotropic matter, a scalar field with an exponential potential and a uniform magnetic field. We shall discuss some general properties of the complete class of models, and then investigate the special case of the Bianchi type-I system in detail. We note that this is not the general class of Bianchi type-I cosmological models; we are only considering those Bianchi type-I models that occur as subcases of the Bianchi type- VI_0 models with one magnetic degree of freedom. It is also of interest to study the Bianchi type-II cosmologies, since although they are very special within the full Bianchi class, they play a central role since the Bianchi type-II state space is part of the boundary of the state space for all higher Bianchi types (i.e., all types except for I and V). The Bianchi type-II models will be investigated in detail elsewhere [31].

The Bianchi type-I system in the absence of a scalar field was first qualitatively analyzed in Ref. [23]. A qualitative analysis of the Bianchi type-I models in the absence of a magnetic field was given in Ref. [20]; in this work the well-known power-law inflationary solution [17,32] was shown to be a stable attractor for an appropriate parameter range in the presence of a barotropic fluid in all Bianchi class B models (previous analysis had shown that this power-law inflationary solution is a global attractor in spatially homogeneous models in the absence of a perfect fluid, except for a subclass of Bianchi type-IX models which recollapse). One of the aims of the present analysis is to study the stability of this model with respect to magnetic field perturbations.

II. THE BIANCHI TYPE- VI_0 MODELS

We shall follow the approach of Ref. [24] in which the theory of dynamical systems was used to give a detailed analysis of the evolution of orthogonal Bianchi cosmologies of type VI_0 with a perfect fluid and a magnetic field as source. This work extended that of Ref. [33] which had studied perfect fluid Bianchi cosmologies of class A using state variables that are dimensionless and have a direct physical or geometric interpretation in terms of the shear of the cosmological fluid, the spatial curvature and the magnetic field, leading to a state space that is a compact subset of \mathbb{R}^5 , which implies that a unified treatment of the asymptotic behavior of the models at early and late times can be given. We note that all of the equilibrium points of the differential equation correspond to self-similar exact solutions of the Einstein field equations.

An invariant orthonormal frame of vector fields on the spacetime are introduced in which one vector is aligned along the fluid flow vector so that the remaining three spatial vectors (triad) span the tangent space orthogonal to the fluid flow at each point of the group orbits. The commutation functions of this frame are then taken as the basic variables. The Einstein-Maxwell field equations for a pure magnetic

field are then derived in the orthonormal frame formalism, giving rise to a set of evolution equations for the shear variables ($\sigma_{\alpha\beta}$), the curvature variables ($n_{\alpha\beta}$), the energy density (conservation equation), the magnetic field (Maxwell equation) and a first integral (the Friedmann constraint), where the remaining freedom in the choice of spatial tetrad was used to simultaneously diagonalize $\sigma_{\alpha\beta}$ and $n_{\alpha\beta}$.

Expansion-normalized variables are then introduced, leading to a reduced set of evolution equations (the dynamical system) for the dimensionless shear variables Σ_{\pm} , the spatial curvature variables N_{α} , the dimensionless magnetic field variables and the density parameter (defined as usual). Defining a dimensionless time variable, the reduced evolution equations in the special case of a Bianchi type-VI₀ were then established (in this case there is a single dimensionless magnetic field variable, M).

A. The equations

We assume a scalar field with an exponential potential $V = V_0 e^{k\phi}$, where V_0 and k are positive constants, and a separately conserved perfect which satisfies the barotropic equation of state $p = (\gamma - 1)\rho$, where the constant γ satisfies $0 \leq \gamma \leq 2$ (although we shall only be interested in the range $0 < \gamma < 2$ here). The energy conservation and Klein-Gordon equations become

$$\dot{\rho} = -3\gamma H\rho, \quad (2.1)$$

$$\dot{\phi} = -3H\dot{\phi} - kV, \quad (2.2)$$

where H is the Hubble parameter, an overdot denotes ordinary differentiation with respect to time t , and units have been chosen so that $8\pi G = 1$.

We also assume a uniform magnetic field (in the “ x ” direction) which satisfies the Maxwell equations. The form of this magnetic field, which we shall denote here as M^2 , is given in Refs. [23] or [24], but its precise form is unimportant since we use the generalized Friedmann constraint to eliminate M^2 from the evolution equations.

We define

$$\Omega \equiv \frac{\rho}{3H^2}, \quad \Phi \equiv \frac{\dot{\phi}}{\sqrt{6}H}, \quad \Psi \equiv \frac{V}{3H^2}, \quad (2.3)$$

two normalized (dimensionless) shear variables Σ_+ and Σ_- (appropriately combinations of the derivatives of the metric functions divided by the Hubble parameter [21]), and the new logarithmic time variable τ by

$$\frac{d\tau}{dt} \equiv H. \quad (2.4)$$

The evolution equations for the quantities

$$X = (\Sigma_+, \Sigma_-, \Phi, \Psi, \Omega, N_+, N_-) \in \mathbb{R}^7 \quad (2.5)$$

are then as follows:

$$\begin{aligned} \Sigma'_+ &= (q-2)\Sigma_+ + 2(1-N_+^2 - \Phi^2 - \Psi - \Sigma_+^2 - \Sigma_-^2 - \Omega) \\ &\quad - 2N_-^2, \end{aligned} \quad (2.6)$$

$$\Sigma'_- = (q-2)\Sigma_- - 2N_+N_-, \quad (2.7)$$

$$\Phi' = (q-2)\Phi - \frac{\sqrt{6}k}{2}\Psi, \quad (2.8)$$

$$\Psi' = (2q+2 + \sqrt{6}k\Phi)\Psi, \quad (2.9)$$

$$\Omega' = (2q-3\gamma+2)\Omega, \quad (2.10)$$

$$N'_+ = (q+2\Sigma_+)N_+ + 6\Sigma_-N_-, \quad (2.11)$$

$$N'_- = (q+2\Sigma_+)N_- + 2\Sigma_-N_+, \quad (2.12)$$

where a prime denotes differentiation with respect to the time τ and the deceleration parameter q is defined by $q \equiv -(1 + H'/H)$, where

$$q = 1 - N_-^2 + \Sigma_+^2 + \Sigma_-^2 + \frac{1}{2}(3\gamma-4)\Omega + \Phi^2 - 2\Psi. \quad (2.13)$$

The physical state space is restricted to $\Omega \geq 0$ and $N_+^2 \leq 3N_-^2$ (where we include equality to include models of Bianchi types I and II [24]). Since $N_- = 0$ implies $N'_- = 0$, the subsets $N_- > 0$ and $N_- < 0$ are invariant. Due to a discrete symmetry in the evolution equations we can, without loss of generality, restrict attention to the subset $N_- \leq 0$.

The magnetic field, defined through the first integral, is given by

$$\frac{3}{2}M^2 = 1 - N_-^2 - \Phi^2 - \Psi - \Sigma_+^2 - \Sigma_-^2 - \Omega, \quad (2.14)$$

and since $M^2 \geq 0$, all of the variables are bounded:

$$0 \leq \{\Sigma_+^2, \Sigma_-^2, \Phi^2, \Psi, \Omega, N_+^2/3, -N_-\} \leq 1. \quad (2.15)$$

Without loss of generality, we can restrict attention to the subset $M \geq 0$. There is also an auxiliary equation for the magnetic field:

$$M' = (q-1-2\Sigma_+)M. \quad (2.16)$$

B. Discussion

A number of invariant sets of the physical state space can be identified (cf., Refs. [24] and [20]). In addition, there are a number of monotonic functions that exist in various invariant sets (particularly on the boundary; cf. p. 521 of Ref. [24]). However, we shall not present them here. Indeed, there are many (over 50) equilibrium points, the vast majority of which are saddles, and hence in this section we shall simply comment on some of the more physically interesting local properties of the models. In the next section we shall investigate the Bianchi I models more rigorously, and a monotonic

function in this invariant set will be given explicitly, in order to discuss some aspects of the intermediate or transient behavior of the models.

The equilibrium point \mathcal{P} : $(\Sigma_+, \Sigma_-, \Phi, \Psi, \Omega, N_+, N_-) = (0, 0, -k/\sqrt{6}, 1 - k^2/6, 0, 0, 0)$, $M=0$ corresponds to the power-law inflationary zero-curvature (flat) FLRW model. The eigenvalues of \mathcal{P} are $\{k^2 - 4, k^2 - 3\gamma, \frac{1}{2}(k^2 - 2)(\times 2), \frac{1}{2}(k^2 - 6)(\times 3)\}$. \mathcal{P} exists for $k^2 \leq 6$ and is a sink for $k^2 < 4$ and $k^2 < 3\gamma$, and is inflationary ($q < 0$) for $k^2 < 2$. In fact it is the global attractor for this range of parameter values; this result generalizes previous work by including a cosmic magnetic field (cf. Ref. [20]).

We are particularly interested in whether there are any sinks (or sources) with $M \neq 0$ in the Bianchi type-VI₀ state space. We know from Ref. [24] that there are no attractors in the absence of a scalar field, so hence we must have $M \neq 0$, $\Phi^2 \neq 0$ (and $\Psi \neq 0$, since this leads to $M=0$). From Eq. (2.10) we see that Ω approaches zero at late times, so we also take $\Omega=0$ (and $q < 2$). From Eqs. (2.8), (2.9) and (2.16) we have that (for $M \neq 0$)

$$\Phi = \frac{-4}{\sqrt{6}k}(1 + \Sigma_+), \quad \Psi = \frac{4}{3k^2}(1 + \Sigma_+)(1 - 2\Sigma_+);$$

$$q = 1 + 2\Sigma_+. \quad (2.17)$$

We are particularly interested in the inflationary case in which $q < 0$; i.e., $-1 < \Sigma_+ < -\frac{1}{2}$. If $N_-, N_+(\Sigma_-)$ are not zero, Eqs. (2.11) and (2.12) yield $N_+ = \pm \sqrt{3}N_-$, $\Sigma_- = \mp (1/2\sqrt{3})(1 + 4\Sigma_+)$, whence Eq. (2.7) then yields $N_-^2 = \frac{1}{12}(1 - 2\Sigma_+)(1 + 4\Sigma_+)$, which is negative for the range of Σ_+ under consideration! [There do exist equilibrium points with $N_- \neq 0$, $N_+ \neq 0$ and $\Sigma_- \neq 0$, but these all have $q \geq 0$ (i.e., are noninflationary) and are *saddles*.] Hence we consider $N_- = N_+ = \Sigma_- = 0$. Equations (2.6) and (2.13) then yield (since $\Sigma_+ \neq \frac{1}{2}$), $\Sigma_+ = -1 + 3k^2/(8 + k^2)$. This equilibrium point, denoted \mathcal{V} in Sec. V, is physical and in the state space when $4 < k^2 < 8$ and is always a saddle (N_+^2 is increasing as orbits evolve away from \mathcal{V} to the future). However, since $q < 0$ only for $k^2 < \frac{8}{3}$ ($\Sigma_+ < -\frac{1}{2}$), this point is never inflationary.

We conclude that inflation is not fundamentally affected by the presence of primordial magnetism in the models under study here.

In Ref. [24] it was shown that in the absence of a scalar field that for $\Omega > 0$, $1 < \gamma < 2$ the global sink is the equilibrium point $PM(VI_0)$ so that at late times all such magnetic Bianchi type-VI₀ cosmologies are approximated by the dynamics of a self-similar Dunn-Tupper [34] magnetovacuum Bianchi type-VI₀ model. This equilibrium point has an analogue here with $\Phi = \Psi = 0$, but is now a saddle. Indeed, it can be easily shown that none of the equilibrium points with vanishing scalar field can be sinks. In particular, the sink in the case $\gamma = 1$ in Ref. [24] [subset of $\mathcal{LM}(VI_0)$] becomes a saddle.

All equilibrium points corresponding to models with a nonzero matter fluid, including those corresponding to the

flat FLRW perfect fluid solution (Einstein–de Sitter) and the flat FLRW scaling solution, are saddles.

In the above we have made some passing comments concerning the intermediate behavior of the models, but we have not discussed this in detail. In the next section we will discuss the intermediate behavior of the (subset) of Bianchi type-I models. The transient behavior is important for assessing the physical significance of the models. For example, Ref. [24] also considered whether any of the Bianchi type-VI₀ magnetic cosmologies are compatible with observations of a highly isotropic universe. For a subset of the Bianchi type-VI₀ cosmologies, the orbits will initially be attracted to the equilibrium point \mathcal{F} corresponding to the flat FLRW model, but will eventually be repelled by it. For such a model there will be an epoch during which the model is close to isotropy and hence will be compatible with observations. This is the phenomenon of *intermediate isotropization*. It will not occur for a randomly selected magnetic Bianchi type-VI₀ model, but there will be a finite probability that an arbitrarily selected model will undergo intermediate isotropization (the probability will depend on the desired closeness to isotropy and the length of the epoch of isotropization) [24].

There is an equilibrium set $\mathcal{K}:(\Sigma_+, \Sigma_-, \Phi, 0, 0, 0, 0)$ where $1 - \Sigma_+^2 - \Sigma_-^2 - \Phi^2 = 0$ (and $M=0$, $q=2$). (These points reduce to the equilibrium set \mathcal{K}_M in the absence of a magnetic field consisting of Jacobs type-I solutions with no matter and a massless scalar field [20]). The eigenvalues are $\{0, 0, 2(1 - 2\Sigma_+), 3(2 - \gamma), 2(1 + \Sigma_+ \pm \sqrt{3}\Sigma_-), 6(1 + k\Phi/\sqrt{6})\}$. There are two zero eigenvalues (corresponding to the fact that \mathcal{K} is a 2-parameter family of equilibrium points). An analysis shows that there is always a subset of \mathcal{K} that acts as sources (these are, in fact, global sources); all others are saddles.

This contrasts with the situation in the absence of a scalar field in which Ref. [24] showed that there are no equilibrium points that are sources and argue that generically an orbit with $\Omega > 0$ will pass through a transient stage and then approach the Kasner circle into the past (consisting of Kasner, Bianchi type-I vacuum models). Since these equilibrium points are saddles, the orbit subsequently leaves along a uniquely determined heteroclinic orbit and then approaches the Kasner circle again. This process then continues indefinitely, and the orbit follows an infinite heteroclinic sequence of Rosen orbits and Taub orbits joining Kasner equilibrium points undergoing Mixmaster-like chaotic oscillations (the past attractor here consists of the union of the Kasner circle and a family of Rosen magnetovacuum type-I orbits and Taub vacuum type-II orbits—see also Ref. [28]).

III. BIANCHI TYPE-I MODELS

The evolution equations for the Bianchi type-I models are obtained by restricting the above equations to the zero-curvature case (i.e., $N_{\pm} = 0$), whence we obtain

$$X = (\Sigma_+, \Sigma_-, \Phi, \Psi, \Omega) \in \mathbb{R}^5, \quad (3.1)$$

and the evolution equations are as follows:

$$\Sigma'_+ = (q-2)\Sigma_+ + 2(1-\Phi^2-\Psi-\Sigma_+^2-\Sigma_-^2-\Omega), \quad (3.2)$$

$$\Sigma'_- = (q-2)\Sigma_-, \quad (3.3)$$

$$\Phi' = (q-2)\Phi - \frac{\sqrt{6}k}{2}\Psi, \quad (3.4)$$

$$\Psi' = (2q+2+\sqrt{6}k\Phi)\Psi, \quad (3.5)$$

$$\Omega' = (2q-3\gamma+2)\Omega, \quad (3.6)$$

where a prime denotes differentiation with respect to the time τ and the deceleration parameter q is now given by

$$q = 1 + \Sigma_+^2 + \Sigma_-^2 + \frac{1}{2}(3\gamma-4)\Omega + \Phi^2 - 2\Psi. \quad (3.7)$$

The magnetic field satisfies

$$\frac{3}{2}M^2 = 1 - \Phi^2 - \Psi - \Sigma_+^2 - \Sigma_-^2 - \Omega, \quad (3.8)$$

and since $M^2 \geq 0$, all of the variables are bounded: $0 \leq \{\Sigma_+^2, \Sigma_-^2, \Phi^2, \Psi, \Omega\} \leq 1$.

Equations (3.1)–(3.6) can also be easily derived from the metric directly in coordinate form—see the Appendix.

A. Invariant sets and monotonic functions

The existence of monotonic functions rules out periodic and recurrent orbits in the phase space and enables global results to be obtained. When $\Omega \neq 0$, we can define $Z \equiv \Omega \Sigma_-^{-2}$, whence from the evolution equations above we find that

$$Z' = 3(2-\gamma)Z, \quad (3.9)$$

so that Z is a monotonically increasing function for $\gamma < 2$ [24]. From the monotonicity principle [21] we can deduce that $\Omega \rightarrow 0$ at early times and $\Sigma_-^2 \rightarrow 0$ at late times.

When $\Omega = 0$, we note that

$$q-2 = 3(M^2 + \Sigma_+^2 + \Sigma_-^2 + \Phi^2 - 1) \leq 0, \quad (3.10)$$

and so from Eq. (3.3) the variable Σ_- is itself a monotonic function.

There are a number of invariant sets. $M=0$ is an invariant set corresponding to the absence of a magnetic field. The dynamics in this invariant set was studied in Ref. [20]. $\Phi = \Psi = 0$ is an invariant set corresponding to the absence of a scalar field. The dynamics in this invariant set was studied in Ref. [23]. $\Omega = 0$ is an invariant set corresponding to the absence of barotropic matter. This invariant set is important for studying the the early time asymptotic behavior. Finally, $\Sigma_- = 0$ is an invariant set which is important in the study of late time asymptotic behavior.

B. Stability of equilibria

Let us present all of the equilibrium points and their associated eigenvalues. We shall assume here that $1 \leq \gamma < 2$. The case of the bifurcation value $\gamma = 2$ will be discussed in Ref. [31]. All equilibrium points correspond to self-similar cosmological models.

Equilibrium points: $\{\Sigma_+, \Sigma_-, \Phi, \Psi, \Omega\}$.

1. Zero magnetic field

\mathcal{P} —Power-law inflation (flat FLRW):

$$\left\{ 0, 0, -\frac{k}{\sqrt{6}}, 1 - \frac{k^2}{6}, 0 \right\}, \quad M=0, \quad q = \frac{1}{2}k^2 - 1.$$

Eigenvalues: $[k^2 - 4, k^2 - 3\gamma, \frac{1}{2}k^2 - 3, \frac{1}{2}k^2 - 3, \frac{1}{2}k^2 - 3]$.

Exists provided $k^2 \leq 6$. Sink if $k^2 < 4$ and $k^2 < 3\gamma$; saddle otherwise. Inflationary if $k^2 < 2$ (note that this is the *only* equilibrium point that can be inflationary).

\mathcal{F}_5 —Matter-scaling (flat FLRW [35]):

$$\left\{ 0, 0, -\frac{\sqrt{6}}{2} \frac{\gamma}{k}, \frac{3}{2} \frac{\gamma}{k^2} (2-\gamma), \frac{k^2-3\gamma}{k^2} \right\}, \quad M=0,$$

$$q = \frac{3}{2}\gamma - 1 > 0.$$

Eigenvalues: $[\frac{3}{2}(\gamma-2), \frac{3}{2}(\gamma-2), 3\gamma-4, \frac{3}{4}(\gamma-2) \pm (3/4k)\sqrt{(2-\gamma)(24\gamma^2+k^2[2-9\gamma])}]$.

Exists provided $k^2 > 3\gamma$. Sink if $\gamma < \frac{4}{3}$; saddle otherwise.

\mathcal{F} —Einstein–de Sitter (flat FLRW):

$$\{0, 0, 0, 0, 1\} \quad M=0, \quad q = \frac{3}{2}\gamma - 1 > 0.$$

Eigenvalues: $[\frac{3}{2}(\gamma-2), \frac{3}{2}(\gamma-2), \frac{3}{2}(\gamma-2), 3\gamma, 3\gamma-4]$.

Saddle.

\mathcal{K} —Kasner vacuum/massless scalar field:

$$\{\Sigma_+^2 + \Sigma_-^2 + \Phi^2 = 1, \quad \Psi = \Omega = 0\}, \quad M=0, \quad q=2.$$

Eigenvalues: $[0, 0, 3(2-\gamma), 2-4\Sigma_+, 6+\sqrt{6}k\Phi]$.

This is a 2-parameter family of equilibrium points. Sources when $\Sigma_+ < \frac{1}{2}$; saddles otherwise.

2. Nonzero magnetic field

These all correspond to anisotropic models ($\Sigma_+^2 + \Sigma_-^2 \neq 0$).

\mathcal{J} —Matter and no scalar field (Jacobs magnetic field model [25]):

$$\left\{ \frac{1}{4}(3\gamma-4), 0, 0, 0, \frac{3}{2} \left(1 - \frac{1}{4}\gamma \right) \right\},$$

$$M^2 = \frac{1}{8}(2-\gamma)(3\gamma-4), \quad q = \frac{3}{2}\gamma - 1 > 0.$$

TABLE I. Summary of all of the equilibrium points (as defined in the text) and their nature for different $\{\gamma, k^2\}$ -parameter values in the Bianchi type-I models with (a) $1 \leq \gamma \leq \frac{4}{3}$ and (b) $\gamma > \frac{4}{3}$. Note that for $\frac{4}{3} < \gamma < 2$, $4 < 3\gamma < 8\gamma(4-\gamma)^{-1} < 8$, and in (b) we have identified $8\gamma(4-\gamma)^{-1} > 6$ (i.e., $\gamma > \frac{12}{7}$) to be specific. For all values of γ , a subset of \mathcal{K} are sources (the remainder are saddles).

(a) $1 \leq \gamma \leq \frac{4}{3}$

k^2	3γ	4	6	8
\mathcal{P}	SINK	SADDLE		DOES NOT EXIST
\mathcal{F}_S	DOES NOT EXIST	SINK		
\mathcal{F}	SADDLE			
\mathcal{V}	DOES NOT EXIST		SADDLE	DOES NOT EXIST

(b) $\frac{4}{3} < \gamma < 2$

k^2	4	3γ	6	$8\gamma/(4-\gamma)$	8
\mathcal{P}	SINK	SADDLE	DOES NOT EXIST		
\mathcal{F}_S	DOES NOT EXIST		SADDLE		
\mathcal{F}	SADDLE				
\mathcal{J}	SADDLE				
\mathcal{S}	DOES NOT EXIST			SINK	
\mathcal{V}	DOES NOT EXIST	SINK	SADDLE	DOES NOT EXIST	

Eigenvalues:

$$\left[\frac{3}{2}(\gamma-2), \frac{3}{2}(\gamma-2), 3\gamma, \frac{3}{4}\{(\gamma-2) \pm \sqrt{(2-\gamma)(3\gamma^2-17\gamma+18)}\} \right].$$

Exists provided $\gamma > \frac{4}{3}$ (same as \mathcal{F} for $\gamma = \frac{4}{3}$). Saddle.

\mathcal{S} —Matter and scalar field:

$$\left\{ \frac{1}{4}(3\gamma-4), 0, -\frac{\sqrt{6}}{2} \frac{\gamma}{k}, \frac{3}{2} \frac{\gamma}{k^2}(2-\gamma), \frac{3}{8k^2}(4k^2-k^2\gamma-8\gamma) \right\},$$

$$M^2 = \frac{1}{8}(2-\gamma)(3\gamma-4), \quad q = \frac{3}{2}\gamma - 1 > 0.$$

Eigenvalues:

$$\left[\frac{3}{2}(\gamma-2), \frac{3}{8}\{2(\gamma-2) \pm (1/k)\sqrt{2(2-\gamma)(a \pm b)}\} \right]$$

where

$$a = 24\gamma^2 + 3\gamma^2k^2 - 26k^2\gamma + 20k^2,$$

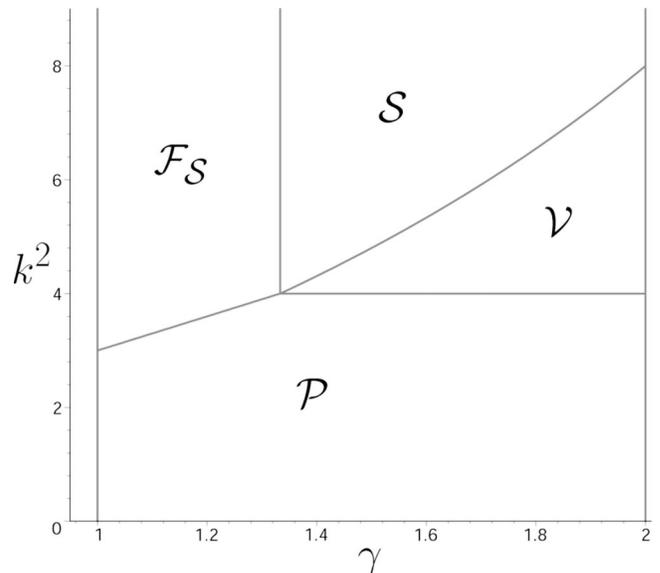


FIG. 1. Summary of all of the sinks for the Bianchi type-I magnetic field models in $\{\gamma, k^2\}$ -parameter space. Note that for each specific pair $\{\gamma, k^2\}$, there is a unique sink.

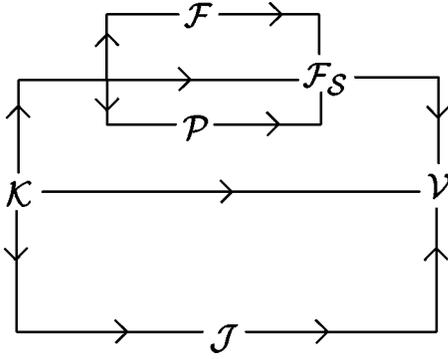


FIG. 2. Heteroclinic sequences for the Bianchi type-I models when $\gamma = \frac{5}{3}$, $5 < k^2 < \frac{40}{7}$.

and

$$b = \sqrt{(18k^2\gamma - 4k^2 + a - 32k\gamma)(18k^2\gamma - 4k^2 + a + 32k\gamma)}.$$

Exists provided $\gamma > \frac{4}{3}$ (same as \mathcal{F}_S for $\gamma = \frac{4}{3}$) and $\gamma \leq 4k^2/(k^2 + 8)$. Sink.

\mathcal{V} —Scalar field and no matter:

$$\left\{ 2\frac{k^2 - 4}{k^2 + 8}, 0, -\frac{2\sqrt{6}k}{k^2 + 8}, 12\frac{8 - k^2}{(k^2 + 8)^2}, 0 \right\},$$

$$M^2 = 2\frac{(k^2 - 4)(8 - k^2)}{(k^2 + 8)}, \quad q = \frac{5k^2 - 8}{k^2 + 8} > 0.$$

Eigenvalues:

$$\begin{aligned} & [3(k^2 - 8)/(k^2 + 8), 3(k^2 - 8)/(k^2 + 8), 3(4k^2 - k^2\gamma) \\ & - 8\gamma)/(k^2 + 8), \frac{3}{2}\{(k^2 - 8)/(k^2 + 8) \\ & \pm \sqrt{(17k^2 - 72)/(k^2 - 8)}\}]. \end{aligned}$$

Exists provided $4 \leq k^2 \leq 8$. Sink when $\gamma > 4k^2/(k^2 + 8)$; saddle otherwise.

C. Discussion

In Table I we summarize the equilibrium points and their stability for different values of the parameters (γ, k^2) . In Fig. 1 we show the sinks in (γ, k^2) -parameter space. The equilibrium point \mathcal{P} represents the power-law FLRW solution with no matter and no magnetic field, and is inflationary for $k^2 < 2$. It is a global attractor for $k^2 < 3\gamma$ in the Bianchi type-I invariant set. Note that \mathcal{P} is a sink in the full class of Bianchi type- VI_0 models, unlike all other sinks in Fig. 1 which are only attractors in the Bianchi type-I invariant set. Also note that the monotonic function Z in the Bianchi type-I invariant set indicates that $\Omega \rightarrow 0$ at early times (at the sources) and $\Sigma_- \rightarrow 0$ at late times (at the sinks), consistent with the results in Table I.

The equilibrium points \mathcal{S} and \mathcal{V} , in which all of the matter fields are nonvanishing and in which matter field (only) vanishes, respectively, correspond to new exact self-similar magnetic field cosmological models. The new solutions are

given explicitly in the Appendix.

Heteroclinic sequences (and hence the transient behavior of the Bianchi type-I models) can easily be deduced from the table. As a particular example, for $\gamma = 5/3$ and $5 < k^2 < 40/7$, a subset of \mathcal{K} constitute sources (the remainder are saddles), \mathcal{V} is a sink, \mathcal{J} is a saddle (that lies in the $\Phi = 0, \Psi = 0$ ($M = 0$) submanifold) and \mathcal{F} , \mathcal{P} and \mathcal{F}_S are saddles (that lie in the isotropic $\Sigma_+ = 0, \Sigma_- = 0$ ($M = 0$) submanifold—in this submanifold \mathcal{K} is a source, \mathcal{F} and \mathcal{P} are saddles, and \mathcal{F}_S is a sink). The heteroclinic sequence in this case is presented in Fig. 2.

From the table and figures we can deduce both the asymptotic and the intermediate behavior of the Bianchi type-I models. Note that as γ increases, the magnetic field becomes increasingly important at late times, consistent with the observation of Wainwright [30] (in the absence of a scalar field).

IV. CONCLUSIONS

We have investigated the class of open Bianchi type- VI_0 universe models with barotropic matter, a scalar field with an exponential potential and a uniform magnetic field. We discussed some of the general properties of the models by utilizing dynamical systems techniques. The equilibrium point \mathcal{P} (with $M = 0$), corresponding to the power-law inflationary flat FLRW model, is a global sink for $k^2 < 2$. Hence all models are future asymptotic to this inflationary attractor for these parameter values. There is an equilibrium set \mathcal{K} with $1 - \Sigma_+^2 - \Sigma_-^2 - \Phi^2 = 0$ (and $M = 0, q = 2$), a subset of which are global sources. Hence all models are past asymptotic to massless scalar field models with no matter and no magnetic field.

A (partial) analysis of the saddles was undertaken in order to determine some of the transient features of the models. In particular, we found that there are no equilibrium points with $M \neq 0$ in the Bianchi type- VI_0 state space which are inflationary, and hence concluded that inflation is not fundamentally affected by the presence of a uniform primordial magnetic field in these models.

This latter result is not necessarily inconsistent with the conclusions of Matravers and Tsagas [8], since only a uniform magnetic field was considered here. Indeed, it is the nonuniform magnetic field gradients that give rise to the magnetocurvature effects in their work which modify the cosmological expansion rate of an almost-FLRW universe and that can have undesirable implications for inflationary models (and may even prevent inflation taking place in the presence of primordial magnetism). It might be thought that on large scales that the magnetic field will be approximately homogeneous, but Matravers and Tsagas [8] have argued that even weak magnetic effects may have significant cosmological consequences. The drawback of their work is that only a local perturbation analysis was performed and questions of genericity and long term behavior cannot be easily addressed. Clearly, an investigation of the qualitative properties of a class of scalar field cosmological models with an inhomogeneous magnetic field would extend and generalize both the work of Matravers and Tsagas [8] and the (restricted) analysis of a uniform magnetic field here.

In order to investigate the intermediate behavior of the models, and hence their physical properties, we discussed the (subset) of Bianchi type-I models in more detail. We found

the equilibrium points \mathcal{S} and \mathcal{V} , in which all of the matter fields are nonvanishing and in which the matter field (only) vanishes, respectively (corresponding to new exact self-similar magnetic field cosmological models), act as sinks in the Bianchi type-I subset. Heteroclinic sequences (and hence the transient behavior of the Bianchi type-I models) were discussed. We found that as γ increases, the magnetic field becomes increasingly important at late times.

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APPENDIX

In the main text we followed the dynamical systems theory approach of LeBlanc *et al.* [24] to analyze the evolution of orthogonal Bianchi cosmologies of type VI₀ with a uniform magnetic field. This approach uses an invariant orthonormal frame in which the commutation functions are the basic variables. Expansion-normalized variables are then introduced, and a reduced set of evolution equations (the dynamical system) for the dimensionless shear variables and spatial curvature variables as well as the dimensionless magnetic field variables and the density parameter are then obtained. All of the equations in the text follow from the analysis of Ref. [24], wherein all quantities are properly defined. However, a coordinate approach is also possible, and it may be of use for the reader to get a better sense of how the equations are derived, at least in the case of the Bianchi type-I models.

The metric for the class of anisotropic Bianchi type-I models, which are the simplest spatially homogeneous generalizations of the flat FLRW models which have nonzero shear but zero three-curvature, is given by

$$ds^2 = -dt^2 + X^2 dx^2 + Y^2 dy^2 + Z^2 dz^2, \quad (\text{A1})$$

where X, Y, Z are functions of t only. The expansion rate is given by

$$\theta \equiv 3H = \frac{\dot{X}}{X} + \frac{\dot{Y}}{Y} + \frac{\dot{Z}}{Z}, \quad (\text{A2})$$

and when $Y = Z$ [the locally rotationally symmetric subcase] there is only one independent rate of shear, which is given by

$$\sigma = \frac{1}{3} \left(\frac{\dot{Y}}{Y} - \frac{\dot{X}}{X} \right). \quad (\text{A3})$$

Normalized (dimensionless) variables and a logarithmic time variable τ are defined by Eqs. (2.3) and (2.4), which then lead to the evolution Eqs. (3.1)–(3.6) in coordinate form [23].

Exact solutions

We can write down the exact solutions corresponding to the equilibrium points \mathcal{S} and \mathcal{V} in the text.

\mathcal{S} : Since $q = \frac{3}{2}\gamma - 1$, we can immediately integrate the Raychaudhuri equation $H'/H = -(1+q)$ to obtain $H = (2/3\gamma)t^{-1}$. Since $\Sigma_- = 0$ and $\Sigma_+ = \frac{1}{4}(3\gamma - 4)$, it follows that $Y = Z$, so that $\sigma = \Sigma_+ \times H = \frac{1}{2}(1 - 4/3\gamma)t^{-1}$. The metric is consequently of power-law form, viz.

$$ds^2 = -dt^2 + t^{2p_1} dx^2 + t^{2p_2} (dy^2 + dz^2), \quad (\text{A4})$$

where p_1 and p_2 are constants. From the expressions for H and σ we then find that $p_1 + 2p_2 = 2/\gamma$ and $p_2 - p_1 = \frac{3}{2} - 2/\gamma$ and hence

$$p_1 = \frac{2}{\gamma} - 1, \quad p_2 = \frac{1}{2}. \quad (\text{A5})$$

The geometry is that of the Jacobs magnetic field model [25], and the only difference here is that the matter fields constituting the source are given by

$$\rho \equiv 3\Omega H^2 = \left(\frac{4k^2 - k^2\gamma - 8\gamma}{2k^2\gamma^2} \right) t^{-2},$$

$$H_1 \equiv 3\sqrt{2}MH = \frac{\sqrt{(2-\gamma)(3\gamma-4)}}{\gamma} t^{-1}, \quad (\text{A6})$$

and the scalar field is given by

$$\phi = \phi_0 - \frac{2}{k} \ln(t), \quad (\text{A7})$$

where $\phi_0 = (1/k) \ln[2(2-\gamma)/\gamma k^2 V_0]$. This solution exists provided $\gamma > \frac{4}{3}$.

\mathcal{V} : In this solution $q = (5k^2 - 8)/(k^2 + 8)$, so that $H = [(k^2 + 8)/6k^2]t^{-1}$. $\Sigma_- = 0$ and $\Sigma_+ = 2(k^2 - 4)/(k^2 + 8)$, and so the metric is of the form of Eq. (A4), where now

$$p_1 = \frac{8 - k^2}{2k^2}, \quad p_2 = \frac{1}{2}. \quad (\text{A8})$$

There is no perfect fluid in this solution ($\rho = 0$), and the magnetic field is given by

$$H_1 = \frac{\sqrt{(k^2 - 4)(8 - k^2)(k^2 + 8)}}{k} t^{-1}. \quad (\text{A9})$$

The scalar field is given by equation (A7) with $\phi_0 = 1/k \ln[(8 - k^2)/(k^4 V_0)]$. This new solution exists provided $4 \leq k^2 \leq 8$.

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